

Appendix A

Unitary decomposition for a one plus two-body Hamiltonian for spinless fermions

Let us consider a system of m fermions in N sp states with a (1+2)-body Hamiltonian $H = h(1) + V(2)$ where $h(1) = \sum_i \epsilon_i \hat{n}_i$ and $V(2)$ is defined by the two-body matrix elements $V_{ijkl} = \langle kl | V(2) | ij \rangle$. With respect to the $U(N)$ group, the two-body interaction $V(2)$ can be separated into scalar ($\nu = 0$), effective one-body ($\nu = 1$) and irreducible two-body ($\nu = 2$) parts. Then, we have [Ch-71, Ko-01a],

$$\begin{aligned}
 V^{\nu=0} &= \frac{\hat{n}(\hat{n}-1)}{2} \bar{V} ; \quad \bar{V} = \binom{N}{2}^{-1} \sum_{i < j} V_{ijij} , \\
 V^{\nu=1} &= \frac{\hat{n}-1}{N-2} \sum_{i,j} \zeta_{i,j} a_i^\dagger a_j ; \quad \zeta_{i,j} = \left[\sum_k V_{kikj} \right] - \left[(N)^{-1} \sum_{r,s} V_{rsrs} \right] \delta_{i,j} , \\
 V^{\nu=2} &= V - V^{\nu=0} - V^{\nu=1} \iff V_{ijkl}^{\nu=2} ; \\
 V_{ijij}^{\nu=2} &= V_{ijij} - \bar{V} - (N-2)^{-1} (\zeta_{i,i} + \zeta_{j,j}) , \\
 V_{ijik}^{\nu=2} &= V_{ijik} - (N-2)^{-1} \zeta_{j,k} \text{ for } j \neq k , \\
 V_{ijkl}^{\nu=2} &= V_{ijkl} \text{ for all other cases .}
 \end{aligned} \tag{A1}$$

Similar to Eq. (A1), the $h(1)$ operator will have $\nu = 0, 1$ parts,

$$\begin{aligned} h^{\nu=0} &= \bar{\epsilon} \hat{n}, \quad \bar{\epsilon} = (N)^{-1} \sum_i \epsilon_i, \\ h^{\nu=1} &= \sum_i \epsilon_i^1 \hat{n}_i, \quad \epsilon_i^1 = \epsilon_i - \bar{\epsilon}. \end{aligned} \tag{A2}$$

Then the propagation equations for the m -particle centroids and variances are,

$$\begin{aligned} E_c(m) &= \langle H \rangle^m = m\bar{\epsilon} + \binom{m}{2} \bar{V}, \\ \sigma^2(m) &= \langle H^2 \rangle^m - [E_c(m)]^2 \\ &= \frac{m(N-m)}{N(N-1)} \sum_{i,j} \left\{ \epsilon_i^1 \delta_{i,j} + \frac{m-1}{N-2} \zeta_{i,j} \right\}^2 \\ &\quad + \frac{m(m-1)(N-m)(N-m-1)}{N(N-1)(N-2)(N-3)} \left\langle \left\langle (V^{\nu=2})^2 \right\rangle \right\rangle^2. \end{aligned} \tag{A3}$$